

Homework #3 Solutions.

①

1. Jackson 3.6	Residues.	l (Å)	l_{eff} (Å)	Z	C
poly-L-alanine	100	3.8	36	0.95	
DNA	1000	3.5	300		

Rms end-to-end distances.

$$\sqrt{R_N^2} = l_{eff} N^{1/2}$$

Now for poly-L-alanine $\frac{l_{eff}}{l} = 9.5$ residues per l_{eff}

$$\sqrt{R_N^2} = \sqrt{\frac{100}{9.5}} 36 \text{ Å} = 117 \text{ Å}.$$

For DNA

$$\frac{l_{eff}}{l} = \frac{300}{3.5} = 86 \text{ residues per } l_{eff}$$

$$\sqrt{R_N^2} = \sqrt{\frac{10^3}{86}} 300 \text{ Å} = 1023 \text{ Å}.$$

If we try the other lengths:

poly-L at 10^3 residues $\sqrt{\frac{10^3}{9.5}} 36 \text{ Å} = 369 \text{ Å}$

DNA at 10^2 residues $\sqrt{\frac{10^2}{86}} 300 \text{ Å} = 323 \text{ Å}.$

2. Jackson 3.9

Modulin $T_m = 350 \text{ K}$ and $\Delta H^\circ = 50 \text{ kcal/mole}.$

Chemically modified so that the 5 valine residues go from 3 rotational states to 1.

(2)

Rotational entropy before modification

$$S_{\text{before}} = k_B \log 3^5 = 5 k_B \log 3$$

$$S_{\text{after}} = k_B \log 1^5 = 0$$

Change in entropy: $-5R \log 3$ per mole.

Change in enthalpy: 0

$$\Delta G = G_{\text{unfolded}} - G_{\text{folded}} = \Delta H^\circ - T_m \Delta S^\circ$$

$$\Rightarrow \Delta H^\circ = H_{\text{unfolded}}^\circ - H_{\text{folded}}^\circ; \quad \Delta S^\circ = S_{\text{unfolded}}^\circ - S_{\text{folded}}^\circ$$

with modification

$$\Delta G^* = \Delta H^\circ - T_m^* \Delta S^{\circ*}$$

$$\Delta S^{\circ*} = S_{\text{unfolded}}^\circ - (S_{\text{folded}}^\circ - 5R \log 3) = \Delta S^\circ + 5R \log 3$$

~~$$\Delta G^* = \Delta H^\circ - T_m^* (\Delta S^\circ + 5R \log 3)$$~~

$$\Delta G^* = \Delta H^\circ - T_m^* (\Delta S^\circ + 5R \log 3) = 0$$

and $\Delta S^\circ = T_m^{-1} \Delta H^\circ$

$$\Delta G^* = \Delta H^\circ - T_m^* \left[\frac{\Delta H^\circ}{T_m} + 5R \log 3 \right] = 0$$

$$\Delta H^\circ \left(1 - \frac{T_m^*}{T_m} \right) = + T_m^* 5R \log 3$$

$$\Delta H^\circ = + T_m^* \left(5R \log 3 + \frac{\Delta H^\circ}{T_m} \right); \quad T_m^* = \frac{\Delta H^\circ}{\frac{\Delta H^\circ}{T_m} + 5R \log 3}$$

Now, we can put in numbers...

③ $R = 1.99 \times 10^{-3} \frac{\text{kcal}}{\text{mole K}}$ so

$$\frac{\Delta H^\circ}{RT_m} = \frac{50}{350 \cdot 1.99 \times 10^{-3}} = 72$$

so $\frac{T_m^*}{T_m} = \frac{72}{72 + 5 \log 3} = 0.93$ or $T_m^* = 326 \text{ K}$.

3. Jackson 4.1

Derive $[P_i] = \frac{[L]^i [P_0]}{k_1 k_2 \dots k_i}$

from $\frac{[P_i]}{[L][P_{i-1}]} = \frac{1}{k_i}$

for $i=1 \Rightarrow [P_1] = \frac{1}{k_1} [L][P_0]$

Assume $[P_i] = \frac{[L]^i [P_0]}{k_1 k_2 \dots k_i}$ and show this holds for

i.

$$[P_i] = \frac{[L][P_{i-1}]}{k_i} = \frac{[L]^i [P_0]}{k_1 \dots k_{i-1} k_i} \quad \checkmark$$

4. Jackson 4.4

$$\Delta G_t^\circ = -RT \log \left(\frac{1 \text{ \AA}^3}{\lambda^3} \right) + RT \log \left(\right)$$

$$R \cong 2 \times 10^{-3} \text{ Kcal/mole} \cdot \text{K} \Rightarrow \frac{\Delta H^\circ}{RT_m} \cong 72$$

$$\Rightarrow \frac{T_m^*}{T_m} = \frac{72}{72 + R \log 3} = 0.93 \Rightarrow T_m^* = 326 \text{ K}$$

3. Jackson 4.1

Derive $[P_i] = [L]^i [P_0] \prod_{j=1}^i K_j$

from $\frac{[P_i]}{[L][P_{i-1}]} = 1/K_i$

By induction: for $i=1$ $[P_1] = \frac{[L][P_0]}{K_1}$

Assume $[P_{i-1}] = \frac{[L]^{i-1} [P_0]}{K_1 K_2 \dots K_{i-1}}$

Then $[P_i] = \frac{[L]^i [P_0]}{(K_1 K_2 \dots K_{i-1}) K_i}$ ✓

4. Jackson 4.4.

Translational partition function $Z = \frac{1}{N!} (V/\lambda^3)^N$ for N noninteracting particles.

$$\Rightarrow G = -k_B T \log \left(\frac{1}{N!} (V/\lambda^3)^N \right) = k_B T (N \log N - N) + k_B T \log (V/\lambda^3)$$

$$= + N k_B T \left(\log \left[\left(\frac{N}{V} \right) \lambda^3 \right] - 1 \right)$$

$$G = N k_B T \log \left(\frac{c \lambda^3}{e} \right) \text{ for one mole } N \rightarrow N_A \text{ of particles.}$$

$$\Rightarrow G^\circ = RT \log \left(c \lambda^3 / e \right)$$

~~$$G^\circ = RT \log \left(\frac{c \lambda^3}{e} \right)$$~~

$$\Delta G_t^\circ = RT \log\left(\frac{1}{(1A)^3} \frac{\lambda^3}{e}\right) - RT \log\left(\frac{N_A}{1L} \frac{\lambda^3}{e}\right) \quad (5)$$

$$\Delta G_t^\circ = \underset{298K.}{RT} \log\left(\frac{(10cm)^3}{N_A (1A)^3}\right) \Rightarrow \frac{(10^{-1}m)^3}{6 \times 10^{23} (10^{-10}m)^3}$$

$$= \frac{1}{6} 10^4$$

$$\Delta G_t^\circ = 4.4 \text{ kcal}$$

5.



$$H = \int_0^L dx \left\{ \frac{1}{2} K (\partial_x^2 u)^2 + \frac{1}{2} \tau (\partial_x u)^2 \right\}$$

Fourier transform: to diagonalize the Hamiltonian

a) Length of the filament:

$$L_0 = \int_0^{L_0} dx \sqrt{1 + (\partial_x u)^2} \approx L_0 + \int_0^{L_0} \frac{1}{2} (\partial_x u)^2 dx$$

↑
small undulation

$$\Rightarrow L_0 - L = \frac{1}{2} \int_0^L (\partial_x u)^2 dx$$

$$b) u(x) = \sum_q u_q \sin(qx) \quad q = \frac{n\pi}{L} \quad n=1,2,\dots$$

$$H = \frac{L}{2} \sum_{n=1}^{\infty} u_n^2 \left[\frac{1}{2} K q_n^4 + \frac{1}{2} \tau q_n^2 \right]$$

Use equipartition theorem

$$\langle U_n \rangle = \frac{2}{L} \frac{k_B T}{K q_n^4 + \tau q_n^2}$$

$$\Rightarrow \langle L \rangle = L_0 - \frac{1}{2} \times 2 \sum_n \frac{k_B T}{K q_n^4 + \tau q_n^2} q_n^2$$

two polarization states.

$$\Rightarrow \langle L \rangle = L_0 - \sum_{n=1}^{\infty} \frac{k_B T}{K q_n^4 + \tau} = L_0 - \frac{k_B T L^2}{\pi^2 K} \sum_{n=1}^{\infty} \frac{1}{\pi n^2 + \tau/K}$$

$$\text{where } \tau_0 = \frac{K \pi^2}{L^2}$$

Now Taylor expand for small τ/τ_0 :

$$\langle L \rangle = L_0 - \frac{k_B T L^2}{\pi^2 K} \sum_{n=1}^{\infty} \frac{1}{n^2} + \frac{k_B T L^2}{\pi^2 K} \frac{\tau}{\tau_0} \sum_{n=1}^{\infty} \frac{1}{n^4} + \dots$$

$$\sum_{n=1}^{\infty} \frac{1}{n^2} = \pi^2/6 \quad \sum_{n=1}^{\infty} \frac{1}{n^4} = \pi^4/90$$

$$\Rightarrow \langle L \rangle = L_0 - \frac{k_B T L^2}{6K} + \frac{k_B T L^4}{K^2 90} \tau + \dots$$

$$d) \Rightarrow \delta L = \frac{k_B T L^4}{K^2 90} \tau \Rightarrow$$

$$K_{\text{eff}} = \frac{90 K^2}{k_B T L^4}$$

6.

Correlation function

$$\langle S_j S_{j+n} \rangle = \frac{1}{Z} \frac{1}{\mathcal{T}} \sum_{\{s_i = \pm 1\}} e^{-\beta H[\{s_i\}]} s_j s_{j+n}$$

Look at the numerator:

$$Z \langle S_j S_{j+n} \rangle = \frac{1}{\mathcal{T}} \sum_{\{s_i = \pm 1\}} \langle s_j | T | s_2 \rangle \cdots \langle s_{j-1} | T | s_j \rangle s_j$$

$$\langle s_j | T | s_{j+1} \rangle \cdots \langle s_{j+n-1} | T | s_{j+n} \rangle s_{j+n} \langle s_{j+n} | T | s_{j+n+1} \rangle \cdots \langle s_N | T | s_1 \rangle$$

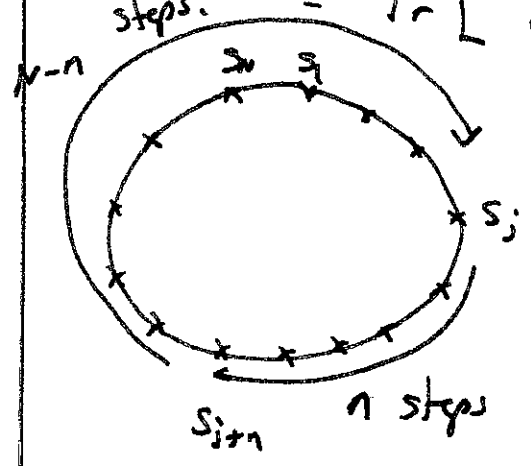
Looking at the far terms corresponding to $s_j = \pm$
 $s_{j+n} = \pm$

We find (as shown in class)

$$Z \langle S_j S_{j+n} \rangle = \text{Tr} [T^{j-1} \sigma_z T^n \sigma_z T^{N-(j+n)+1}]$$

Pauli Spin matrix $\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$$= \text{Tr} [T^{N-n} \sigma_z T^n \sigma_z]$$



Now we diagonalize T ($\omega/h=0$) as in class (8)
and plug in:

$$\sum \langle S_j S_{j+n} \rangle = \text{Tr} \left[A \begin{pmatrix} \lambda_1^{N-n} & 0 \\ 0 & \lambda_2^{N-n} \end{pmatrix} A \begin{pmatrix} \lambda_1^n & 0 \\ 0 & \lambda_2^n \end{pmatrix} \right]$$

where $A = \begin{pmatrix} \Delta & \Delta + \xi \\ -\Delta + \xi & -\Delta \end{pmatrix} \frac{1}{\xi} = \begin{pmatrix} 0 & 1 \\ +1 & 0 \end{pmatrix}$

and $\Delta = \frac{J}{2}$ and $\xi = 2 e^{-\beta J}$
for $h=0$.

Complete the algebra, you get (see class notes for more

$$C(n) = \langle S_j S_{j+n} \rangle = \left(\frac{e^{\beta J} - e^{-\beta J}}{e^{\beta J} + e^{-\beta J}} \right)^n \stackrel{\text{details}}{=} \tanh^n(\beta J)$$

exponential decay of correlation.

$$C(n) = \exp[n \log(\tanh(\beta J))]$$

c)

or we have a correlation length of

$$n^* = \frac{1}{|\log[\tanh(\beta J)]|} \text{ sites. } e^{-n/n^*} = C(n)$$

